## Nuclear Quadrupole Relaxation and PCl<sub>3</sub> Group Reorientations in Crystals

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The nuclear quadrupole relaxation time  $T_1$  temperature dependence is studied experimentally for the <sup>35</sup>Cl nuclei in some Cl<sub>3</sub>P==NR-type crystals. The conventional additive libration-reorientation scheme of the relaxation mechanism is accepted. Using this scheme, the  $T_1$  temperature dependence experimental data are analyzed. The results obtained make it possible to prove the existence of the PCl<sub>3</sub> group autonomous reorientations and to find the values of all the parameters characterizing the relaxation processes.

It is well known that different kinds of inner motions are of great importance to the nuclear quadrupole relaxation processes in molecular crystals, which makes it possible to study these motions by means of nuclear quadrupole relaxation measurements. The mobility of the CCl<sub>3</sub> group, possessing a threefold symmetry axis, has recently been widely investigated by the nuclear quadrupole relaxation method (1, 2). The PCl<sub>3</sub> group, which has the same symmetry, has not been studied so widely. A peculiar feature of the reorientation motions of the PCl<sub>3</sub> group is that they occur with respect to the direction of a "formally double" bond (as in the Cl<sub>3</sub>P==NR type compounds, for example). The PCl<sub>3</sub> group motions in the Cl<sub>3</sub>P==NSO<sub>2</sub>Cl crystal have been investigated (3) using NQR line width measurements. In this paper, the inner motions in some Cl<sub>3</sub>P==NR-type crystals are studied using  $T_1$  temperature dependence measurements for <sup>35</sup>Cl nuclei. The mobility types are identified and the numerical values of the corresponding parameters are found.

The experimental curves  $T_1(T)$  have been obtained using the progressive pulse saturation method (4); here T is temperature in the interval between 77 K and the melting point. The curves having no phase transition singularities, it is natural to accept the well-known additive libration-reorientation scheme of the relaxation mechanism (1, 5):

$$(T_1^{-1})_{obs} = (T_1^{-1})_{libr} + (T_1^{-1})_{reor} = a T^n + b \exp\left(-\frac{V_0}{RT}\right).$$
 [1]

To get the parameters  $a, n, b, V_0$ , the experimental data have been fitted by means of the conventional least-squares method using computer techniques. Two procedures are possible: one can use the curve Eq. [1] as a whole (1), or get the term  $(T_1^{-1})_{reor}$  as the difference between  $(T_1^{-1})_{obs}$  and  $(T_1^{-1})_{libr}$ , the latter being obtained from the low-temperature data (4). Both procedures give practically the same values of the parameters,

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TABLE 1

		PCI <sub>3</sub> group resonance frequencies	F		Paramet	ers of Eq. (1)			
Compound	Melting point temperature (K)	at 77 K (MHz)	at 77 K (msec)	a (sec <sup>-1</sup> deg <sup>-n</sup> )	u	$b (\sec^{-1})$	V <sub>o</sub> (kcal/mole)	7, (K)	$\lambda_{f}$ (sec <sup>-1</sup> )
NCCI(CCI <sub>3</sub> )2	326	29.442 29.481 29.764	51 42 62	$\begin{array}{c} 3.31 \times 10^{-3} \\ 4.68 \times 10^{-3} \\ 4.78 \times 10^{-3} \end{array}$	1.99 $1.95$ $1.86$	- 3.10 × 10 <sup>11</sup>	$6.8 \pm 0.2$	180	$2 \times 10^{3}$
NPO(CCl <sub>3</sub> ) <sub>2</sub>	395	29.718 30.070 30.320	79 79 105	1.12 × 10 <sup>-4</sup> 1.12 × 10 <sup>-4</sup> 1.50 × 10 <sup>-4</sup>	2.66 2.66 2.53	· 1.35 × 10 <sup>11</sup>	<b>5.1</b> ± <b>0.1</b>	145	$2 \times 10^{3}$
NCOCF <sub>3</sub>	245	29.600 30.085 30.762	68 96 88	8.32 × 10 <sup>-4</sup> 3.39 × 10 <sup>-4</sup> 1.23 × 10 <sup>-3</sup>	2.24 $2.37$ $2.10$	· 2.00 × 10 <sup>11</sup>	<b>6.2</b> ± <b>0.2</b>	175	$3 \times 10^{3}$
-NSO2N—PCI3		29.302 29.500 30.468	52 44	$5.75 \times 10^{-3}$	06.1	$.5.10 \times 10^{12}$	$6.8 \pm 0.1$	184	3 × 10 <sup>4</sup>
	314	29.355 29.526 30.897	104 101 153	$1.90 \times 10^{-3}$	1.87	$6.61 \times 10^{12}$	$6.6 \pm 0.1$	177	$4 \times 10^4$
-NSO <sub>2</sub> CI	308	30.133 30.254 31.780	161 163 192	$1.82 \times 10^{-3}$	1.83	- 5.75 × 10 <sup>11</sup>	<b>7.3 ± 0.1</b>	225	$4 \times 10^{4}$
-NCCl <sub>3</sub> CCl <sub>3</sub>	295	29.075 29.645 31.152	190 245 192	$8.01 \times 10^{-4}$ $6.65 \times 10^{-4}$ $8.01 \times 10^{-4}$	2.00 2.00 2.00	$2.82 \times 10^9$ $1.24 \times 10^9$ $2.82 \times 10^9$	$3.4 \pm 0.2$ $3.4 \pm 0.2$ $3.4 \pm 0.2$ $3.4 \pm 0.2$	134 151 134	

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but the second one makes it possible to transform the functions  $(T_1^{-1})_{\text{libr}} = \phi(T)$  and  $(T_1^{-1})_{\text{reor}} = \psi(T^{-1})$  into linear ones, which is profitable from the point of view of the error calculation. In this paper both procedures have been utilized.

Table 1 gives the values of the parameters a, n, b,  $V_0$  obtained in the way described above together with some other data of importance: melting points, NQR frequencies, and  $T_1$  values at 77 K, and fade out temperatures  $T_f$  and corresponding reorientation rates estimated as

$$\lambda = \frac{3}{4} (T_1^{-1})_{\text{reor}}$$
<sup>[2]</sup>

[see (1)]. To check the correctness of the values obtained, in some cases the curve for Eq. [1] has been calculated using these values and compared with the experimental curve (see Figs. 1 and 3).



FIG. 1. Temperature dependence of the relaxation time for the <sup>35</sup>Cl nuclei in different molecular fragments in Cl<sub>3</sub>P=NCCl(CCl<sub>3</sub>)<sub>2</sub> (two low-frequency lines of the PCl<sub>3</sub> group coalesce at high temperatures). The solid lines represent the calculated curves with the parameter values from Table 1.  $\circ$ , PCl<sub>3</sub>;  $\bullet$ , CCl<sub>3</sub>;  $\times$ , CCl.

The values of the parameters a, n, b,  $V_0$  provide us with important quantitative characteristics of the inner motions in the crystals considered. In addition to that, some interesting qualitative conclusions concerning these motions may be drawn from the results obtained.

First of all, the existence of the PCl<sub>3</sub> group autonomous reorientations may be taken for granted in some crystals. In fact, according to Fig. 1, showing the  $T_1(T)$  curve for Cl<sub>3</sub>P=NCCCl(CCl<sub>3</sub>)<sub>2</sub> there are two circumstances: (1) In the temperature interval where the curves corresponding to the <sup>35</sup>Cl nuclei in the PCl<sub>3</sub> group become practically exponential, those corresponding to the <sup>35</sup>Cl nuclei in the remainder of the molecule are far from being like that; and (2) in the temperature interval mentioned the PCl<sub>3</sub> group curves are very near to one another. No doubt the first circumstance gives an argument for the existence of separate reorientations of the PCl<sub>3</sub> group. The second one does also if the possible reorientation axes of the molecule as a whole do not coincide with the symmetry axes of the PCl<sub>3</sub> group; this coincidence does not take place for the substances considered. In the Cl<sub>3</sub>P=NPO(CCl<sub>3</sub>)<sub>2</sub> crystal the situation is the same. In the Cl<sub>3</sub>P=NCOCF<sub>3</sub>, Cl<sub>3</sub>P=NSO<sub>2</sub>N=PCl<sub>3</sub>, and Cl<sub>3</sub>P=NSO<sub>2</sub>Cl cases only the second one of the two circumstances mentioned above takes place (for Cl<sub>3</sub>P=NSO<sub>2</sub>N=PCl<sub>3</sub> see Fig. 2). In  $Cl_3P$ =MSO<sub>2</sub>N=PCl<sub>3</sub> and  $Cl_3P$ =MSO<sub>2</sub>Cl, this is the only argument for the existence of the PCl<sub>3</sub> group autonomous reorientations: The point is that in the  $Cl_3P$ =MSO<sub>2</sub>N=PCl<sub>3</sub> molecule all the <sup>35</sup>Cl nuclei are in the PCl<sub>3</sub> groups, and in the  $Cl_3P$ =MSO<sub>2</sub>Cl case the NQR signal from the S-Cl fragment fades out before that from



FIG. 2. Temperature dependence of the resonance frequencies (a) and relaxation times (b) for the <sup>35</sup>Cl nuclei in  $Cl_3P$ —NSO<sub>2</sub>N—PCl<sub>3</sub>. Two sets of curves: 1, 3, 5 and 2, 4, 6 concern two PCl<sub>3</sub> groups in the same molecule. Relaxation data are given only for the upper frequency lines corresponding to each of the two sets.



FIG. 3. Temperature dependence of the relaxation time for the <sup>35</sup>Cl nuclei in different molecular fragments in Cl<sub>3</sub>P=NCCl<sub>2</sub>CCl<sub>3</sub> (the PCl<sub>3</sub> solid lines represent the calculated curves with the parameter values from Table 1).  $\circ$ , PCl<sub>3</sub>;  $\bullet$ , CCl<sub>3</sub>;  $\times$ , CCl<sub>2</sub>.

the PCl<sub>3</sub> group (3). In the Cl<sub>3</sub>P=NCOCF<sub>3</sub> case there is a complementary argument given by some <sup>19</sup>F NMR data (6).

On the other hand, situations are possible where the separate reorientations of the PCl<sub>3</sub> group cannot be observed, because the NQR signal disappears early due to the reorientations of the molecules as a whole. For example, this occurs in the Cl<sub>3</sub>P== NCCl<sub>2</sub>CCl<sub>3</sub> crystal (see Fig. 3) where the  $T_1(T)$  curves for different molecular fragments give the same reorientation potential barrier values.

In the cases where the separate reorientations of the  $PCl_3$  group have been found, the corresponding barrier values lie in the interval 5–7 kcal/mole (see Table 1). Taking into account the fact that the barrier in question is due to the inner interactions in molecules on the one hand and to the intermolecular interactions on the other hand, one sees that the hindering potential for the  $PCl_3$  group reorientations due to the formally double bond P=N is rather low.

The nuclear quadrupole relaxation time  $T_1$  is highly sensitive to the reorientation motions; for example, in the case of the <sup>35</sup>Cl nuclei in the PCl<sub>3</sub> group the essential shortening of  $T_1$  begins at reorientation rate values from the interval 2–10 sec<sup>-1</sup>, see Fig. 1 and 2b.

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